Star Formation Efficiency in Nearby Galaxies Revealed with a New CO-to-H₂ Conversion Factor Prescription

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Submitted to ApJL

ABSTRACT

Determining how galactic environment, especially the high gas densities and complex dynamics in bar-fed galaxy centers, alters the star formation efficiency (SFE) of molecular gas is critical to understanding galaxy evolution. However, these same physical or dynamical effects also alter the emissivity properties of CO, leading to variations in the CO-to-H₂ conversion factor ($\alpha_{\rm CO}$) that impact the assessment of the gas column densities and thus of the SFE. To address such issues, we investigate the dependence of $\alpha_{\rm CO}$ on local CO velocity dispersion at 150 pc scales using a new set of dust-based

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 $\alpha_{\rm CO}$ measurements, and propose a new $\alpha_{\rm CO}$ prescription that accounts for CO emissivity variations across galaxies. Based on this prescription, we estimate the SFE in a sample of 65 galaxies from the PHANGS-ALMA survey. We find increasing SFE towards high surface density regions like galaxy centers, while using a constant or metallicity-based $\alpha_{\rm CO}$ results in a more homogeneous SFE throughout the centers and disks. Our prescription further reveals a mean molecular gas depletion time of 700 Myr in the centers of barred galaxies, which is overall 3–4 times shorter than in non-barred galaxy centers or the disks. Across the galaxy disks, the depletion time is consistently around 2–3 Gyr regardless of the choice of $\alpha_{\rm CO}$ prescription. All together, our results suggest that the high level of star formation activity in barred centers is not simply due to an increased amount of molecular gas but also an enhanced SFE compared to non-barred centers or disk regions.

Keywords: CO line emission (262) — Disk galaxies (391) — Galaxy nuclei (609) — Giant molecular clouds (653) — Star formation (1569)

1. INTRODUCTION

Star formation in galaxies is governed by the amount of molecular gas and the efficiency with which that gas is converted into stars. To understand the evolutionary process of star formation activity within galaxies, it is critical to measure the molecular gas star formation efficiency (SFE; defined as the ratio between star formation rate, SFR, and molecular gas mass, M_{mol}), or molecular gas depletion time ($t_{\rm dep} = 1/{\rm SFE}$) (see review by Saintonge & Catinella 2022). Previous studies have found that SFR and molecular gas surface densities are highly correlated (i.e., the molecular Kennicutt-Schmidt relation, or mKS relation; Kennicutt 1998) and that $t_{\rm dep}$ is usually at 1–4 Gyr across nearby star-forming galaxies (e.g., Bigiel et al. 2008; Leroy et al. 2008; Saintonge et al. 2011; Schruba et al. 2011; Utomo et al. 2017; Sun et al. 2023). Despite the minor variation in general, $t_{\rm dep}$ is also found to vary systematically with local and global host galaxy properties, which could be driven by environmental and/or dynamical effects from e.g., metallicity, molecular cloud structure, bar instabilities, active galactic nuclei, or galaxy interactions (Saintonge et al. 2011, 2012; Schruba et al. 2019; Ellison et al. 2021a,b; Querejeta et al. 2021; Lu et al. 2022; Villanueva et al. 2022; Jiménez-Donaire et al. 2023; Maeda et al. 2023).

The assessment of molecular gas SFE relies heavily on the CO-to-H₂ conversion factor $(\alpha_{CO})^1$:

$$\alpha_{\rm CO} = \frac{M_{\rm mol}}{L'_{\rm CO(1-0)}} = \frac{\Sigma_{\rm mol}}{I_{\rm CO(1-0)}} \left[\frac{\rm M_{\odot}}{\rm K~km~s^{-1}~pc^2} \right] , (1)$$

where M_{mol} (Σ_{mol}) is the total molecular gas mass (surface density) and $L'_{\text{CO}(1-0)}$ ($I_{\text{CO}(1-0)}$) is the line lumi-

nosity (intensity) of CO J=1-0. $\alpha_{\rm CO}$ is known to vary with molecular gas conditions such as density, temperature, and dynamical state (see review by Bolatto et al. 2013, hereafter B13), which are the same conditions that could also alter the intrinsic SFE of the molecular gas. Due to the lack of a widely-agreed prescription that can accurately predict $\alpha_{\rm CO}$, many studies could only assume a constant $\alpha_{\rm CO}$ referencing the Milky Way (MW) disk average (e.g., B13) to convert CO observations to molecular gas mass. This has made $\alpha_{\rm CO}$ variation one of the dominant sources of uncertainty in current molecular gas and SFE studies (see discussions in Ellison et al. 2020a; Maeda et al. 2023; Sun et al. 2023).

The impacts of $\alpha_{\rm CO}$ variations on both SFE and cloud evolutionary timescale estimates are particularly critical in galaxy centers (Leroy et al. 2013; Utomo et al. 2017; Muraoka et al. 2019; Ellison et al. 2020a; Pessa et al. 2021; Maeda et al. 2023; Sun et al. 2023). In those environments, $\alpha_{\rm CO}$ can be 5–15 times lower than the Galactic disk value (Ackermann et al. 2012; Sandstrom et al. 2013; Israel 2020; Teng et al. 2022, 2023; den Brok et al. 2023). The lower $\alpha_{\rm CO}$ in galaxy centers is likely driven by CO emissivity variations due to higher excitation and/or stronger dynamical effects such as turbulence or inflowing gas (Narayanan et al. 2012; Papadopoulos et al. 2012; Bolatto et al. 2013; Gong et al. 2020; Teng et al. 2023). These effects may also explain the low $\alpha_{\rm CO}$ seen in mergers or (ultra-)luminous infrared galaxies (U/LIRGs) (Downes & Solomon 1998; Krieger et al. 2017; Sliwa et al. 2017; Cicone et al. 2018; Herrero-Illana et al. 2019).

Reducing the uncertainty in molecular gas and SFE studies, and thereby improving our understanding in star formation and galaxy evolution, requires a robust $\alpha_{\rm CO}$ prescription that can be systematically applied to large samples of galaxies with diverse environments. Recent studies have proposed various types of $\alpha_{\rm CO}$ pre-

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 $^{^1}$ $\alpha_{\rm CO}$ is defined for the CO J=1-0 line in most literature, but it can also be evaluated for other transitions. In this work, when we refer to $\alpha_{\rm CO}$, we mean $\alpha_{\rm CO(1-0)}$ unless otherwise specified.

scription depending on metallicity, stellar mass surface density, SFR, SFE, and/or CO line-related properties (Genzel et al. 2012; Narayanan et al. 2012; Bolatto et al. 2013; Hunt et al. 2015; Amorín et al. 2016; Accurso et al. 2017; Renaud et al. 2019; Gong et al. 2020; Madden et al. 2020; Ramambason et al. 2023). However, establishing a reliable $\alpha_{\rm CO}$ calibration remains a challenge because it requires $\alpha_{\rm CO}$ measurements covering a sufficient sample of galaxies spanning a broad range of molecular gas physical and dynamical conditions, and the two most realistic ways to measure $\alpha_{\rm CO}$ in nearby galaxies are via dust emission (which is typically restricted to kpc resolutions; Israel 1997; Leroy et al. 2011; Sandstrom et al. 2013; Schruba et al. 2017; den Brok et al. 2023; Yasuda et al. 2023; Chiang et al. in prep) or multi-CO isotopologue observations (which is expensive at cloud scales; Sliwa et al. 2017; Cormier et al. 2018; Israel 2020; Teng et al. 2022, 2023).

Thanks to the high resolution and sensitivity of the Atacama Large Millimeter/submillimeter Array (ALMA), CO (isotopologue) observations are now routinely possible at cloud scales in nearby galaxies (e.g., Leroy et al. 2021; Davis et al. 2022; Koda et al. 2023; Williams et al. 2023). In particular, recent studies modeling multi-CO isotopologues in nearby galaxy centers have revealed that CO opacity is the dominant driver of $\alpha_{\rm CO}$ variations (Israel 2020; Teng et al. 2022, 2023). This strong dependence of $\alpha_{\rm CO}$ on CO opacity further leads to a clear anti-correlation between $\alpha_{\rm CO}$ and the observed line width at \sim 100 pc scales in barred galaxy centers (Teng et al. 2023, hereafter T23).

Motivated by these latest measurements of $\alpha_{\rm CO}$, we will test if the correlation found in T23 also applies to the 12 galaxies (labeled with * in Table 1, including 8 barred and 4 non-barred) which have dust-inferred $\alpha_{\rm CO}$ values at kpc scales (from Chiang et al. in prep) and molecular gas velocity dispersion measured at 150-pc scales (from the PHANGS–ALMA survey; Leroy et al. 2021; Sun et al. 2022). The results of this comparison

lead us to a new $\alpha_{\rm CO}$ prescription capturing CO emissivity effects in star-forming galaxies. In this Letter, we present this prescription, discuss its physical implications, and study its impact on SFE across a sizable sample of galaxy centers and disks with diverse properties.

2. DATA AND MEASUREMENTS

2.1. PHANGS Datasets

Our analysis is based on various molecular gas and star formation properties, leveraging a database developed by Sun et al. (2022) which assembled multiwavelength measurements of 80 galaxies from the PHANGS-ALMA survey (Leroy et al. 2021). From this database, we extract multiple physical quantities in matched hexagonal apertures with fixed sizes of 1.5 kpc. The quantities used in this work include: intensity-weighted mean molecular gas velocity dispersion measured at 150-pc scale ($\langle \Delta v \rangle_{150\text{pc}}$), area-weighted mean CO(2-1) line integrated intensity $(I_{\text{CO}(2-1)})$, stellar mass surface density (Σ_{star}) , SFR surface density (Σ_{SFR}), and gas-phase metallicity (Z', normalized to the solar value $[12 + \log(O/H)_{\odot} = 8.69]$ and calibrated based on Pettini & Pagel 2004). All these quantities are corrected for the effects of galaxy inclination and data sensitivity limits (see Sun et al. 2022 for more details).

To further explore trends in galaxies with or without stellar bars, we adopt the classification of stellar bars for PHANGS galaxies (Querejeta et al. 2021). Table 1 lists the 65 galaxies included in our analysis, which is the overlap between Querejeta et al. (2021) and Sun et al. (2022). This sample from PHANGS has high-resolution CO(2–1) data with beam sizes of 150 pc or smaller. Columns (6–10) in Table 1 show the measurements extracted from Sun et al. (2022) for the central 1.5 kpc regions of those galaxies.

Table 1. Galaxy Sample and Properties in the Central 1.5 kpc Regions

Galaxy	Bar	Dist.	Incl.	P.A.	Z'	$\log(\Sigma_{\mathrm{SFR}})$	$I_{\text{CO}(2-1)}$	$\log(\Sigma_{\mathrm{star}})$	$\langle \Delta v \rangle_{150 \mathrm{pc}}$	$\log(\alpha_{\mathrm{CO}}^{\mathrm{Eq.2}})$	$\log(t_{\rm dep})$
		$[\mathrm{Mpc}]$	$[\deg]$	$[\deg]$	$[\mathrm{Z}_{\odot}]$	$\left[rac{ m M_{\odot}}{ m yr~kpc^2} ight]$	$[{\rm K~km~s^{-1}}]$	$[{\rm M}_{\odot}~{\rm pc}^{-2}]$	$[{\rm km~s^{-1}}]$	$\left[\frac{\rm M_{\odot}\ s}{\rm K\ km\ pc^2}\right]$	[yr]
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)
IC1954	1	12.8	57.1	63.4	1.10	-1.67	7.5	2.51	7.1	0.36	9.08
IC5273	1	14.2	52.0	234.1	1.12	-1.59	4.7	2.48	7.4	0.34	8.78
NGC0253*	1	3.7	75.0	52.5	1.33	0.3166	198.2	990.2	28.5	-0.15	8.88
NGC0628*	0	9.8	8.9	20.7	1.29	-1.82	6.3	3.03	5.4	0.45	9.26
NGC0685	1	19.9	23.0	100.9	1.23	-2.15	2.8	2.42	6.8	0.37	9.15

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Galaxy	Bar	Dist. [Mpc]	[deg]	P.A. [deg]	Z' $[\mathrm{Z}_{\odot}]$	$\log(\Sigma_{\rm SFR}) \atop \left[\frac{\rm M_{\odot}}{\rm yr~kpc^2}\right]$	$I_{\rm CO(2-1)}$ [K km s ⁻¹]	$\log(\Sigma_{\rm star})$ $[{\rm M}_{\odot}~{\rm pc}^{-2}]$	$\langle \Delta v \rangle_{150 \text{pc}}$ [km s ⁻¹]	$\frac{\log(\alpha_{\rm CO}^{\rm Eq.2})}{\left[\frac{\rm M_{\odot}~s}{\rm K~km~pc^2}\right]}$	$\log(t_{\mathrm{dep}})$ [yr]
(1)	(2)	(3)	(4)	(5)	(6)	[yr kpc²] (7)	(8)	(9)	(10)	[K km pc²] (11)	(12)
NGC1087	1	15.8	42.9	359.1	1.19	-1.12	24.7	2.68	13.4	0.12	8.82
NGC1097	1	13.6	48.6	122.4	1.34	-0.29	196.0	3.61	33.1	-0.20	8.57
NGC1300	1	19.0	31.8	278.0	1.33	-1.57	44.6	3.25	23.2	-0.07	9.33
NGC1317	1	19.1	23.2	221.5	1.33	-1.45	25.2	3.64	18.0	0.02	9.06
NGC1365	1	19.6	55.4	201.1	1.36	0.04	462.4	3.80	25.6	-0.11	8.70
NGC1385	0	17.2	44.0	181.3	1.21	-1.00	15.0	2.78	9.5	0.25	8.60
NGC1433	1	18.6	28.6	199.7	1.35	-1.46	44.6	3.63	17.3	0.03	9.33
NGC1511	0	15.3	72.7	297.0	1.18	-1.17	12.4	2.60	7.5	0.33	8.79
NGC1512	1	18.8	42.5	261.9	1.34	-1.36	31.9	3.41	11.3	0.19	9.24
NGC1546	0	17.7	70.3	147.8	1.29	-1.52	45.0	3.08	7.9	0.31	9.67
NGC1559	1	19.4	65.4	244.5	1.29	-1.47	10.5	2.77	6.6	0.38	9.06
NGC1566	1	17.7	29.5	214.7	1.34	-0.97	55.7	3.63	28.7	-0.15	8.76
NGC1637	1	11.7	31.1	20.6	1.20	-0.82	12.2	2.57	22.8	-0.07	8.03
NGC1792	0	16.2	65.1	318.9	1.33	-1.35	52.0	3.08	10.7	0.21	9.46
NGC1809	0	20.0	57.6	138.2	1.14	-1.71	2.2	2.19	3.1	0.65	8.88
NGC2090	0	11.8	64.5	192.5	1.22	-2.19	4.5	2.79	4.2	0.54	9.57
NGC2283	1	13.7	43.7	-4.1	1.18	-1.93	3.0	2.31	5.1	0.47	9.07
NGC2566	1	23.4	48.5	312.0	1.34	-0.05	265.0	3.16	26.5	-0.12	8.53
NGC2835	1	12.2	41.3	1.0	1.21	-2.23	2.2	$\frac{3.10}{2.47}$	4.3	0.53	9.30
NGC2903	1	10.0	66.8	203.7	1.33	-0.81	55.8	3.15	19.7	-0.01	8.73
NGC2907	0	14.1	33.0	108.1	1.34	-0.95	68.5	3.13	16.2	0.06	9.03
NGC3059	1	20.2	29.4	-14.8	1.29	-0.84	23.1	2.68	11.4	0.18	8.58
NGC3039 NGC3137	0	16.4	70.3	-0.3	1.18	-2.53	3.0	2.19	3.2	0.18	9.85
NGC3351*	1	10.4	45.1	193.2	1.13	-2.55 -0.71	34.4	3.18	19.2	-0.01	8.43
NGC3507	1	23.6	21.7	55.8	1.30	-0.71	17.3	3.16	23.4	-0.01	8.91
NGC3507 NGC3511	1	13.9	75.1	256.8	1.22	-1.94	17.3	2.50	6.6	0.38	9.74
NGC3511*	0	13.9 13.2	68.8	343.0	1.36	-1.66	14.8	3.50	6.7	0.38	9.40
NGC3521	0	13.2 11.3	25.1	78.4	1.10	-1.55	6.3	$\frac{3.50}{2.68}$	7.0	0.36	8.90
NGC3621	0	7.1	65.8	343.8	1.10	-1.86	8.4	2.66	5.2	0.46	9.44
NGC3626	1	20.0	46.6	165.2	1.31	-1.40	11.2	3.52	13.5	0.40	8.76
NGC3627*				173.1		-1.40 -1.19	64.9	$\frac{3.52}{3.53}$	34.6	-0.22	8.98
NGC4254*	$\frac{1}{0}$	11.3 13.1	57.3 34.4	68.1	1.35 1.30	-1.19	35.7	3.27	10.3	0.22	9.06
NGC4294 NGC4293		15.1	65.0	48.3	1.31	-1.10	45.0	2.92	27.4	-0.13	8.82
	1								9.7	0.13	9.38
NGC4298 NGC4303	0	14.9	59.2	313.9 312.4	1.22	-1.79	$14.3 \\ 70.8$	$2.69 \\ 3.52$	9.7 16.9	0.24	9.38 8.91
NGC4303 NGC4321*	1	17.0 15.2	$23.5 \\ 38.5$	312.4 156.2	1.32 1.34	-0.83 -0.78		$\frac{3.52}{3.35}$	16.9 19.5	-0.01	8.91 8.96
	1						101.0				
NGC4457 NGC4496A	1	$15.1 \\ 14.9$	$17.4 \\ 53.8$	$78.7 \\ 51.1$	1.30 1.04	-1.15 2.21	38.7 1.6	3.69 2.17	29.0 2.8	-0.15 0.69	8.77 9.29
NGC4496A NGC4535	1					-2.21					
NGC4535 NGC4536*	1	15.8	44.7 66.0	179.7	1.32	-1.03 0.56	43.1	2.78	20.1	-0.02	8.83 8.76
	1	16.2	66.0	305.6	1.30	-0.56	110.3	3.30	21.0	-0.04	8.76
NGC4540	1	15.8	28.7	12.8	1.14	-1.90	3.9	2.65	6.3	0.40	9.07
NGC4548	1	16.2	38.3	138.0	1.34	-1.96	9.7	3.38	24.8	-0.10	9.03
NGC4569*	1	15.8	70.0	18.0	1.35	-1.13	112.1	3.24	27.6	-0.14	9.23
NGC4571	0	14.9	32.7	217.5	1.23	-2.36	2.5	2.63	2.7	0.70	9.63
NGC4689*	0	15.0	38.7	164.1	1.26	-1.81	9.2	2.61	5.6	0.44	9.40

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Galaxy	Bar	Dist.	Incl.	P.A.	Z'	$\log(\Sigma_{ m SFR})$	$I_{\text{CO}(2-1)}$	$\log(\Sigma_{\mathrm{star}})$	$\langle \Delta v \rangle_{150 \mathrm{pc}}$	$\log(\alpha_{\mathrm{CO}}^{\mathrm{Eq.2}})$	$\log(t_{\mathrm{dep}})$
		[Mpc]	$[\deg]$	$[\deg]$	$[\mathrm{Z}_{\odot}]$	$\left[\frac{\mathrm{M}_{\odot}}{\mathrm{yr}\ \mathrm{kpc}^{2}}\right]$	$[{\rm K~km~s^{-1}}]$	$[{\rm M}_{\odot}~{\rm pc}^{-2}]$	$[{\rm km~s^{-1}}]$	$\left[\frac{\rm M_{\odot}\ s}{\rm K\ km\ pc^2}\right]$	[yr]
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)
NGC4731	1	13.3	64.0	255.4	1.02	-1.97	1.6	1.99	5.1	0.47	8.85
NGC4781	1	11.3	59.0	290.0	1.09	-1.76	9.1	2.59	8.2	0.30	9.20
NGC4826	0	4.4	59.1	293.6	1.27	-1.43	27.9	3.26	21.8	-0.05	9.01
NGC4941*	1	15.0	53.4	202.2	1.25	-1.31	6.9	3.00	18.7	0.00	8.34
NGC4951	0	15.0	70.2	91.2	1.15	-1.77	10.0	2.67	11.4	0.18	9.14
NGC5042	0	16.8	49.4	190.6	1.18	-2.15	2.2	2.52	5.3	0.46	9.14
NGC5068	1	5.2	35.7	342.4	0.98	-2.01	1.0	2.33	3.6	0.60	8.79
NGC5128	0	3.7	45.3	32.2	1.36	-0.78	45.3	3.70	22.7	-0.07	8.56
NGC5134	1	19.9	22.7	311.6	1.30	-1.94	4.9	3.33	11.4	0.18	9.00
NGC5248*	1	14.9	47.4	109.2	1.30	-0.97	72.6	3.35	15.7	0.07	9.08
NGC5530	0	12.3	61.9	305.4	1.23	-2.02	6.4	3.00	4.8	0.50	9.51
NGC5643	1	12.7	29.9	318.7	1.29	-0.29	42.0	2.88	26.7	-0.12	7.98
NGC6300	1	11.6	49.6	105.4	1.31	-0.55	41.1	2.85	36.2	-0.23	8.11
NGC7456	0	15.7	67.3	16.0	1.09	-2.70	0.9	1.99	2.7	0.70	9.56
NGC7496	1	18.7	35.9	193.7	1.21	-0.43	46.2	2.72	23.7	-0.08	8.20

Note. (1) Galaxies with an asterisk are those with $\alpha_{\rm CO}$ measurements (see Section 2.2); (2) bar classification (Querejeta et al. 2021); (3) distance (Anand et al. 2021); (4–5) inclination and position angles (Lang et al. 2020); (6–10) the central 1.5 kpc measurements of gas-phase metallicity (PP04-based), kpc-averaged SFR surface density, kpc-averaged CO(2-1) integrated intensity, kpc-averaged stellar mass surface density, and CO intensity-weighted mean velocity dispersion at 150-pc scale (Sun et al. 2022); (11) $\log(\alpha_{\rm CO})$ derived from (10) using Equation 2; (12) molecular gas depletion time derived from (7), (8), and (11) using Equation 6.

2.2. Dust-based α_{CO} Measurements

We obtain spatially resolved $\alpha_{\rm CO}$ from Chiang et al. (in prep), where $\alpha_{\rm CO}$ is measured at 2 kpc resolution across 41 nearby (≤ 20 Mpc) and moderately-inclined (Incl. $\leq 80^{\circ}$) spiral galaxies with resolved measurements of CO integrated intensity (including PHANGS-ALMA) and atomic gas. The authors assumed a constant dust-to-metals ratio to constrain the total gas mass with dust and metallicity measurements. In their sample, 8 barred and 4 non-barred galaxies from PHANGS have dust-based $\alpha_{\rm CO}$ measurements (those with an * in Table 1). These measurements typically cover out to a galacto-centric radius of ~ 10 kpc, including ~ 2000 Nyquist-sampled data points. It is based on these data that we examine scaling relations of $\alpha_{\rm CO}$ and develop an $\alpha_{\rm CO}$ prescription in Section 3.1.

The $\alpha_{\rm CO}$ measurements by Chiang et al. were derived based on the PHANGS CO(2–1) data, and we directly use their $\alpha_{\rm CO(2-1)}$ measurements to ensure methodological consistency when we derive molecular gas surface density and SFE (see Sections 3.3 and 3.4). To compare with most $\alpha_{\rm CO}$ literature using $\alpha_{\rm CO(1-0)}$, however, we convert the measured $\alpha_{\rm CO(2-1)}$ to $\alpha_{\rm CO(1-0)}$ by assuming a CO(2–1)/(1–0) ratio (R_{21}) of 0.65. Such results can be easily reverted to $\alpha_{\rm CO(2-1)}$ via a linear scaling with

0.65. We note that Chiang et al. also provided $\alpha_{\text{CO}(1-0)}$ measurements assuming a SFR-dependent R_{21} , and we have checked that using such α_{CO} does not change any of our results qualitatively (see Section 3.1).

We also note that the metallicity adopted by Chiang et al. for computing $\alpha_{\rm CO}$ is based on the S-calibration in Pilyugin & Grebel (2016, hereafter PG16S), which is different from the O3N2 calibration used for the PHANGS dataset based on Pettini & Pagel (2004, hereafter PP04). Recent studies suggest that PG16S is a more reliable metallicity prescription than PP04 (e.g., Kreckel et al. 2019). With the data on 12 galaxies, we find that PP04 estimates result in ~ 0.2 dex higher Z' than PG16S (see also De Vis et al. 2019), which might be due to the mismatch in the adopted solar oxygen abundance value under different calibration schemes (e.g., $12 + \log(O/H)_{\odot} = 8.50$ or 8.69; see discussion in Esteban et al. 2022). Throughout this work, we adopt PG16S-based Z' from Chiang et al. (in prep) for analyses restricted to these 12 galaxies. However, due to the lack of PG16S-based measurements on all 65 PHANGS galaxies, we use the PP04-based Z' when implementing metallicity-dependent $\alpha_{\rm CO}$ prescriptions across the full sample for consistency.

To evaluate the credibility of the observed α_{CO} trends with our parameters of interest (i.e., $\langle \Delta v \rangle_{150pc}$ and Z'), we calculate for each parameter bin the number of pixels with reliable $\alpha_{\rm CO}$ measurements divided by the number of pixels with measured Δv or Z'. For Δv , we find the fraction of reliable pixels to be 70-100% for bins with $\langle \Delta v \rangle_{150 \mathrm{pc}} \gtrsim 3 \mathrm{~km~s^{-1}}$, while it drops significantly to < 50% in lower velocity dispersion bins². This means that our $\alpha_{\rm CO}$ data coverage is insufficient to accurately represent regions with $\Delta v \lesssim 3 \text{ km s}^{-1}$. As for Z', the corresponding completeness of $\alpha_{\rm CO}$ is above 60% across regions with $Z' \gtrsim 0.6$, while it drops below 40% at lower metallicities (where the PHANGS-ALMA dataset has poorer coverage). These "incomplete" regimes will be excluded by our fitting and analysis in Section 3.1, where we present the new $\alpha_{\rm CO}$ prescription.

3. RESULTS

3.1. A Velocity Dispersion-based α_{CO} Prescription

To investigate how the dust-based $\alpha_{\rm CO}$ varies with local velocity dispersion, we use nearest-neighbor matching to relate the α_{CO} measurements at 2-kpc scales with the velocity dispersion which is measured at 150pc scale and then averaged over 1.5-kpc-sized apertures via intensity weighting $(\langle \Delta v \rangle_{150pc})$. As shown in Figure 1(a), the data clearly follow an inverse power-law relation, which is in close agreement with the fit by T23 on three barred galaxy centers at ~100 pc scales (i.e., dash-dotted orange line, assuming a CO/H₂ abundance of 1.5×10^{-4}). The central regions of the 12 galaxies (vertical, gray bars in Figure 1(a)) align well with the overall trend, showing that velocity dispersion can trace $\alpha_{\rm CO}$ variations in both the centers and disks³. The green cross sign marks the typical MW disk values of $\alpha_{\rm CO} \sim 4.35 \; \rm M_{\odot} \; (K \; km \; s^{-1} \; pc^2)^{-1} \; and \; \Delta v = 5 \; km \; s^{-1},$ which also agrees with the overall trend. The shaded area in Figure 1(a) indicates the regime where $\alpha_{\rm CO}$ data is incomplete (see Section 2.2).

Excluding the incomplete regime, we conduct a least-squares fitting in log-log space based on the remaining ~1600 data points, using the curve_fit function in scipy.optimize. The best-fit power-law relation to these data from 12 galaxies is represented by

$$\log \alpha_{\rm CO} = -0.81 \, \log \langle \Delta v \rangle_{150 \rm pc} + 1.05 \,, \tag{2}$$

where $\alpha_{\rm CO}$ and $\langle \Delta v \rangle_{150 \rm pc}$ are in units of ${\rm M}_{\odot}$ (K km s⁻¹ pc²)⁻¹ and km s⁻¹, respectively. The best-fit relation is shown by the black dashed line in the top panel of Figure 1(a) and is consistent with the trend of the binned $\alpha_{\rm CO}$ medians. The dispersion of data with respect to Equation 2 is $\sigma \sim 0.12$ dex, and the standard deviation error returned by curve_fit is ± 0.02 for both the fitted slope and intercept. We remind readers that the $\alpha_{\rm CO}$ data here are converted from $\alpha_{\rm CO(2-1)}$ assuming $R_{21}=0.65$, and thus it should be scaled by $R_{21}/0.65$ if R_{21} is known. If a SFR-dependent R_{21} is used following Chiang et al. (in prep), the trend of $\alpha_{\rm CO}$ in Figure 1(a) could be shallower by 30–40%, as indicated by the red dashed line.

While the functional fit in Equation 2 is based on Δv measured at 150-pc scale, we also find a similar best-fit relation (dotted line) for six of those galaxies where $\langle \Delta v \rangle_{\rm 90pc}$ is available. Because Δv does not vary strongly between 90 and 150 pc scales (see also Sun et al. 2022), we would not expect this to change our results, and thus Equation 2 should be applicable with Δv measurements around 100-pc resolutions. We note that the evaluation of Δv can also be affected by the number of gas components overlapping along the same sightlines, which could increase Δv in barred galaxy centers. However, such effect is found to be mild (see T23, Appendix A), and we expect it to be even milder in our case, as $\langle \Delta v \rangle_{150pc}$ is averaged over kpc-sized regions.

3.2. Comparison to Previous Literature

We compare our Δv -based prescription with existing $\alpha_{\rm CO}$ prescriptions in the literature, including those based on metallicity (Accurso et al. 2017; Sun et al. 2020a) or combining metallicity and stellar mass surface density (B13). First, we investigate if metallicity alone could trace the observed $\alpha_{\rm CO}$ variations. Figure 1(b) relates the measured $\alpha_{\rm CO}$ with metallicity, using the same metallicity as those used in Chiang et al. (in prep) to calculate $\alpha_{\rm CO}$ (see Section 2.2). The data and the power-law fit (black dashed line) overall agrees with the purely metallicity-dependent $\alpha_{\rm CO}$ prescription from Accurso et al. (2017) (orange dashed line)⁴, although the data scatter is larger than the trend with velocity dispersion. In the regime where our dataset is complete, the scatter of the observed $\alpha_{\rm CO}$ is $\sigma \sim 0.1$ dex with $\langle \Delta v \rangle_{150 \text{pc}}$ and 0.3 dex with Z'. This shows a significant improvement in predicting $\alpha_{\rm CO}$ with our Δv -based

 $^{^2}$ This is likely due to a large amount of low S/N measurements clustering around ${\sim}2.5$ km/s, which is the velocity resolution of the PHANGS CO data

 $^{^3}$ In this paper, center refers to the central ${\sim}2$ kpc-sized aperture at $R_{\rm gal}{=}0$, and disks represents the rest of the measurements.

⁴ The Z' in the original prescription [$\alpha_{\rm CO} = 4.35(Z')^{-1.6}$] was based on the PP04 calibration. Here we convert their prescription to the same (PG16S-based) metallicity scale as we adopt, using an approximate conversion based on De Vis et al. (2019).

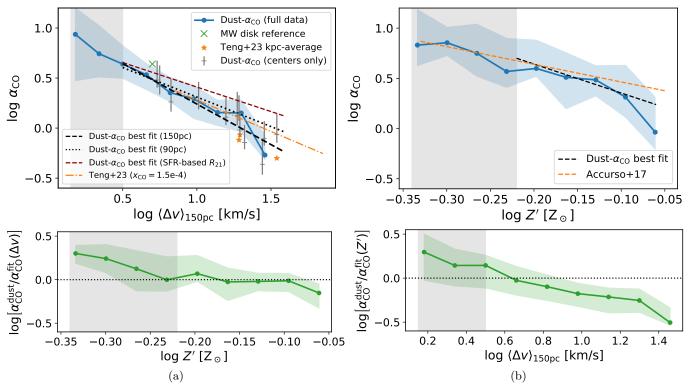


Figure 1. Column (a): Dust-based α_{CO} measurements show a strong anti-correlation with the intensity-weighted average of 150-pc scale molecular gas velocity dispersion (top), consistent with the result from T23 on barred galaxy centers (orange line, with an assumed CO/H₂ abundance $x_{CO} = 1.5 \times 10^{-4}$); the blue lines and shaded area represent the binned medians and 16th–84th percentile of the measured α_{CO} ; the gray shaded area indicates low-confidence regime where α_{CO} sampling is incomplete; the black dashed/dotted lines show the best fit power-law relations with $\langle \Delta v \rangle$ at 150/90 pc resolutions, and the red dashed line represents the best fit relation when α_{CO} is derived by assuming a SFR-dependent R_{21} (Chiang et al. in prep). The residuals of the fit (bottom) do not correlate with Z' in the data-complete regime, suggesting that the observed α_{CO} variations can be fully captured by our Δv -based prescription, without requiring an additional metallicity dependence. Column (b): Similar to (a), but the measured α_{CO} is correlated with metallicity (top), and the residuals are correlated with $\langle \Delta v \rangle_{150pc}$ (bottom); the orange dashed line marks the prediction from Accurso et al. (2017), which agrees with the overall data but shows a larger scatter.

prescription, compared to current metallicity-dependent prescriptions.

In the bottom panels of Figures 1(a) and (b), we relate the residuals of each $\alpha_{\rm CO}$ fit with Z' or $\langle \Delta v \rangle_{150 \rm pc}$, in order to check if metallicity effects can explain any residual variation of $\alpha_{\rm CO}$ around the Δv trend, or the opposite. Above the completeness limit, we find no trend between the residuals from the Δv prescription and metallicity. On the other hand, the residuals from the metallicity fit clearly decrease with Δv above the completeness threshold. This suggests that Δv is crucial for tracing the $\alpha_{\rm CO}$ changes, even without including metallicity effects. We have checked that the $\alpha_{\rm CO}$ correlation with Z' seen in this regime may come from the correlation between Z' and Δv , as both variables decrease with the galactocentric radius.

Taking both metallicity and emissivity effects into account, B13 also suggested a tentative prescription⁵ based on α_{CO} measurements in nearby disks and (U)LIRGs:

$$\alpha_{\rm CO} \approx 2.9 \, \exp\left(\frac{0.4}{Z'}\right) \left(\frac{\Sigma_{\rm star} + \Sigma_{\rm mol}}{100 \, {\rm M}_{\odot} \, {\rm pc}^{-2}}\right)^{-\gamma} ,$$
 (3)

where $\gamma = 0.5$ if $\Sigma_{\rm star} + \Sigma_{\rm mol} > 100~{\rm M}_{\odot}~{\rm pc}^{-2}$ or $\gamma = 0$ otherwise. To compare the derived $\alpha_{\rm CO}$ from our proposed prescription (Equation 2) with that from B13, we apply both prescriptions to galaxies in the PHANGS sample (see Table 1) using kpc-scale Z' and $\Sigma_{\rm star}$. As we find $\Sigma_{\rm mol} \ll \Sigma_{\rm star}$ even with a (likely-overestimated) Galactic $\alpha_{\rm CO}$ (Sun et al. 2022), we neglect $\Sigma_{\rm mol}$ in Equa-

⁵ The original prescription included a molecular cloud surface density term which was assumed at 100 M_{\odot} pc⁻². Here we adopt the same value and note that this helps avoid unrealistic α_{CO} values in low surface density regions (Sun et al. 2023, T23).

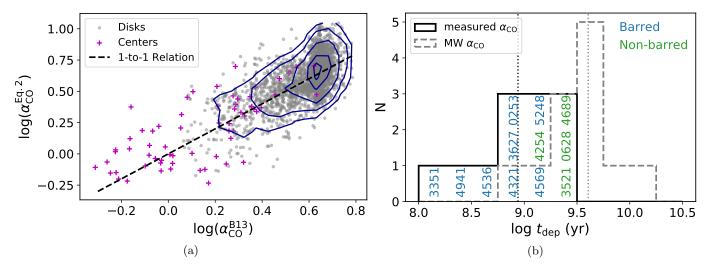


Figure 2. (a) Comparison of the derived $\alpha_{\rm CO}$ using our Δv -based prescription (Equation 2) and the Z' plus $\Sigma_{\rm star}$ -based prescription (Equation 3; B13), applied to 65 galaxies. The overlaid contours indicate 16%, 50%, 84%, 95%, and 98% data inclusion of the disk regions. The two prescriptions show a general 1-to-1 agreement (dashed line), which supports the credibility of our prescription. (b) Molecular gas depletion time ($t_{\rm dep}$) of 12 galaxy centers (with their NGC names shown on the histogram) determined by the measured $\alpha_{\rm CO}$ (solid line) and the MW $\alpha_{\rm CO}$ (dashed line). The median $t_{\rm dep}$ using measured or MW $\alpha_{\rm CO}$ are indicated by the vertical dotted lines. Overall, $t_{\rm dep}$ is lower using the measured $\alpha_{\rm CO}$, and a clear separation is found between barred and non-barred galaxies, suggesting high star formation efficiency in barred galaxy centers.

tion 3. We note that Chiang et al. (in prep) also reported a similar $\alpha_{\rm CO}$ relation that scales with $\Sigma_{\rm star}^{-0.5}$.

Figure 2(a) compares the $\alpha_{\rm CO}$ values predicted by Equations 2 and 3. Excluding the regime of $\log(\alpha_{CO}) \gtrsim$ 0.65 where B13 enforces a MW-like $\alpha_{\rm CO}$ value with $\gamma = 0$ (which also corresponds to the low-confidence regime of our Δv -based prescription), the two prescriptions show an overall match with a ~ 0.5 dex scatter. Despite a significant scatter, this general agreement may indicate that Δv and $\Sigma_{\rm star}$ are tracing the same physical process that drives $\alpha_{\rm CO}$ variations. A likely scenario is that Δv is set by the additional gravitational potential from stellar components, which can thus be tracked by $\Sigma_{\rm star}$ (see B13 and Chiang et al. in prep). It is also possible that Δv is a proxy of molecular gas surface densities and/or local CO intensities which could also reflect opacity and $\alpha_{\rm CO}$ changes, as previous studies have found good correlations between these properties (Sun et al. 2022, see also Section 4 for further discussion).

The B13 prescription was mostly based on $\alpha_{\rm CO}$ measurements that were independent from ours and included several U/LIRGs in their sample, and the ~ 0.5 dex scatter with our prescription is also consistent with the uncertainty estimated by B13. Therefore, the rough agreement seen in Figure 2(a) may also provide additional evidence for the validity of our proposed prescription. Compared to a $\Sigma_{\rm star}$ -based prescription, one advantage of using a Δv -based prescription is that Δv straightforwardly traces the optical depth change (Teng et al. 2022, 2023), making it closer to the underlying physics

that could control $\alpha_{\rm CO}$ variations. Another advantage is that Δv can be directly obtained from the CO data. Thus, no ancillary multi-band data are needed to estimate $\alpha_{\rm CO}$, which circumvents uncertainties in translating observations into $\Sigma_{\rm star}$.

We note that the scaling of $\alpha_{\rm CO}$ with Δv in Equation 2 is similar to what would be predicted by simple theoretical arguments. As shown by Equation 8 in Gong et al. (2020), the excitation temperature $(T_{\rm ex})$ under Large Velocity Gradient approximation with assumptions of a two-level optically-thick system can be written as

$$T_{\rm ex} \propto \rho_{\rm mol} \sqrt{\frac{L_{\rm mol} \cdot x_{\rm CO}}{\Delta v}} ,$$
 (4)

where $\rho_{\rm mol}$ and $L_{\rm mol}$ are the density and size of a COemitting molecular cloud, respectively. To first order, we also have $I_{\rm CO} \sim T_{\rm ex} \cdot \Delta v$ from the cloud. Thus, combining Equation 1 with Equation 4, we obtain

$$\alpha_{\rm CO} = \frac{\Sigma_{\rm mol}}{I_{\rm CO}} \sim \frac{\rho_{\rm mol} \cdot L_{\rm mol}}{T_{\rm ex} \cdot \Delta v} \propto \sqrt{\frac{L_{\rm mol}}{x_{\rm CO} \cdot \Delta v}} \ .$$
 (5)

The resulting $\alpha_{\rm CO}$ dependence on the inverse square root of Δv is similar to the fits in Figure 1. While the fitted slope for $\langle \Delta v \rangle_{150 \rm pc}$ (Equation 2) is slightly steeper than -0.5, we emphasize that the above calculation is highly simplified and is only for providing an intuitive check with theoretical expectations.

3.3. Star Formation Efficiency in Galaxy Centers

As $\alpha_{\rm CO}$ determines the total molecular gas surface density ($\Sigma_{\rm mol}$, in units of ${\rm M}_{\odot}~{\rm pc}^{-2}$), the variation of

 $\alpha_{\rm CO}$ directly affects the estimation of molecular gas depletion time $(t_{\rm dep})$ or SFE $(=1/t_{\rm dep})$:

$$t_{\rm dep} = \Sigma_{\rm mol} / \Sigma_{\rm SFR} = \alpha_{\rm CO} \cdot I_{\rm CO} / \Sigma_{\rm SFR}$$
 . (6)

While we examine only the SFE in this work, we note that the impact of $\alpha_{\rm CO}$ on estimating the SFE per molecular cloud free-fall time is even more significant, as $\alpha_{\rm CO}$ also affects the assessment of cloud density which changes the free-fall time (e.g., Querejeta et al. 2023; Sun et al. 2023). Motivated by the clear trend of galaxy centers having lower $\alpha_{\rm CO}$ values (Figures 1(a) and 2(a)), we derive $t_{\rm dep}$ for the 12 galaxy centers with $\alpha_{\rm CO}$ measurements (Chiang et al. in prep; T23), using kpc-scale $\Sigma_{\rm SFR}$ and $I_{\rm CO(2-1)}$ (see Table 1). Then, we examine how $t_{\rm dep}$ in galaxy centers derived from the measured $\alpha_{\rm CO}$ would differ from that using the standard MW $\alpha_{\rm CO}$ of 4.35 (or 6.7 in terms of $\alpha_{\rm CO(2-1)})$ ${\rm M}_{\odot}$ (K km s $^{-1}$ pc 2) $^{-1}$.

Figure 2(b) presents histograms of $t_{\rm dep}$ for the 12 galaxy centers. For the histogram using the measured $\alpha_{\rm CO}$, we separate barred and non-barred galaxies with different colors. We find that the median $t_{\rm dep}$ with the MW $\alpha_{\rm CO}$ is 4–5 times longer than that with the measured $\alpha_{\rm CO}$. Furthermore, adopting the MW $\alpha_{\rm CO}$ results in a similar $t_{\rm dep}$ of ~ 3 Gyr between barred and non-barred centers. In contrast, if the measured $\alpha_{\rm CO}$ is used, the median $t_{\rm dep}$ of barred and non-barred centers becomes 0.6 Gyr and 2.0 Gyr, respectively, differing by more than a factor of three. This suggests that SFE in barred galaxy centers tend to be higher than non-barred galaxy centers, and that using a constant $\alpha_{\rm CO}$ can obscure such a trend.

3.4. Systematic Impact on Star Formation Efficiency

With the PHANGS sample (Table 1), we investigate the impact of different $\alpha_{\rm CO}$ prescriptions on SFE or $t_{\rm dep}$ in the centers and disks of barred and non-barred galaxies. Figure 3 shows the molecular Kennicutt-Schmidt (mKS) relation across all 65 galaxies measured at the 1.5-kpc scale, comparing $\Sigma_{\rm mol}$ determined from our $\alpha_{\rm CO}$ prescription (Equation 2) with that determined using a MW $\alpha_{\rm CO}$. It is clear that adopting the MW $\alpha_{\rm CO}$ results in a wider range of $\Sigma_{\rm mol}$ with values reaching $> 1000 \text{ M}_{\odot} \text{ pc}^{-2}$ in galaxy centers, while our prescription suggests $\Sigma_{\rm mol} < 200~{\rm M}_{\odot}~{\rm pc}^{-2}$ in general. Furthermore, our prescription reveals a trend of higher SFE towards higher $\Sigma_{\rm mol}$, which steepens the mKS relation for galaxy centers and other high- $\Sigma_{\rm mol}$ regions. With the MW α_{CO} , however, both galaxy centers and disks exhibit a roughly constant SFE. These results show that $\alpha_{\rm CO}$ and $\Sigma_{\rm mol}$ in galaxy centers may overall be overestimated by a factor of 5 with the MW α_{CO} , and that

the choice of $\alpha_{\rm CO}$ greatly affects our understanding of galactic-scale star formation.

Figure 4 presents histograms of velocity dispersion and $t_{\rm dep}$ across the PHANGS sample, separating centers (upper panels) and disks (lower panels) for barred (blue) and non-barred (green) galaxies. In non-barred galaxy centers, $\langle \Delta v \rangle_{150 {\rm pc}}$ is typically $< 10~{\rm km~s^{-1}}$, while barred centers span a significantly wider range up to $\sim 40~{\rm km~s^{-1}}$. On the other hand, barred and non-barred disks show consistent velocity dispersion, with $\langle \Delta v \rangle_{150 {\rm pc}}$ typically below 5 km s⁻¹ but reaching up to 10 km s⁻¹. These distributions agree with Sun et al. (2020b), who reported similar Δv between galaxy disks and non-barred centers but an overall ~ 5 times higher Δv in barred centers.

We then examine the distribution of t_{dep} derived with different $\alpha_{\rm CO}$ prescriptions. Using our Δv -based prescription, we find distinctly different t_{dep} between barred and non-barred centers, with the mean/median of $t_{\rm dep}$ in barred galaxy centers ($\sim 700 \text{ Myr}$) being 3 times shorter than in non-barred centers (~ 2.1 Gyr). The 16th–84th percentile ranges for t_{dep} in barred and non-barred centers is 0.3–1.6 and 0.8–3.6 Gyr, respectively. In contrast, all other prescriptions result in < 0.2 dex difference between the median $t_{\rm dep}$ of the two types of systems. Such a small difference between barred and non-barred centers is even true for the B13 prescription which shows similarly short t_{dep} for all galaxy centers that generally matches our results. Particularly, the MW α_{CO} leads to two completely overlapping $t_{\rm dep}$ distributions, overestimating the overall t_{dep} in barred galaxy centers by a factor of 3-4 if compared to our results. As for the disks, the median of $t_{\rm dep}$ remains consistent at 2–3 Gyr across all four prescriptions, while it is found to be systematically lower in barred galaxies than in non-barred galaxies by ~ 0.1 dex.

Notably, our prescription reveals short $t_{\rm dep}$ down to $\lesssim 100$ Myr in some barred galaxy centers, which is not seen with other prescriptions. Such a short time scale is supported by recent simulations of galaxy centers including effects from bar-driven inflows (e.g., Armillotta et al. 2019; Sormani et al. 2020; Moon et al. 2021). In addition, we note that the overall $t_{\rm dep}$ for galaxy centers is similar between our result and B13's, both suggesting $t_{\rm dep} \sim 1$ Gyr which is shorter than the disks value of ~ 3 Gyr. This factor-of-three difference between centers and disks is consistent with recent simulations (e.g., Tress et al. 2020a). However, using the MW or Z-based $\alpha_{\rm CO}$ for galaxy centers obscures such difference and leads to similar $t_{\rm dep}$ across entire galaxies.

4. DISCUSSION

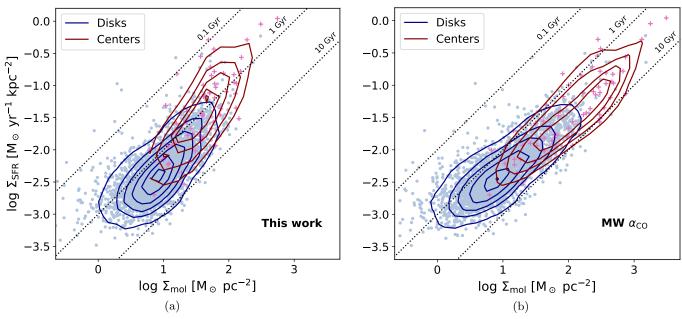


Figure 3. The molecular Kennicutt-Schmidt (mKS) relation across 65 PHANGS galaxies, where the $\alpha_{\rm CO}$ used to derive $\Sigma_{\rm mol}$ is based on (a) Equation 2 or (b) the MW value. The thin dotted lines represent constant molecular gas depletion times ($t_{\rm dep}$) of 0.1, 1, and 10 Gyr. With our $\alpha_{\rm CO}$ prescription, the galaxy centers clearly show a steeper trend than the disks, indicating shorter $t_{\rm dep}$ towards higher $\Sigma_{\rm mol}$. Adopting the MW $\alpha_{\rm CO}$ instead results in a roughly constant $t_{\rm dep}$ for both centers and disks.

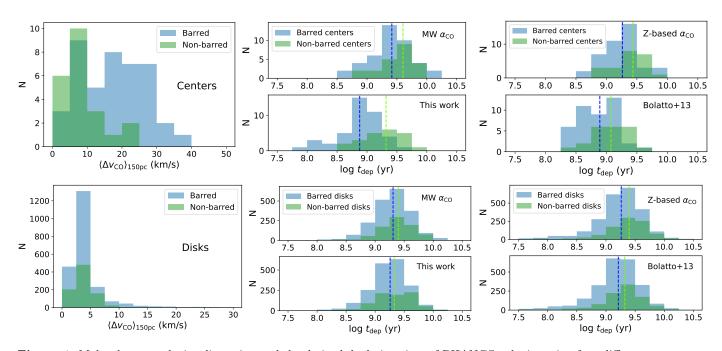


Figure 4. Molecular gas velocity dispersion and the derived depletion time of PHANGS galaxies using four different $\alpha_{\rm CO}$ prescriptions. The upper/lower panels show the centers/disks regions. The medians of the barred/non-barred distributions are indicated by the blue/green dashed lines. Our prescription reveals that barred centers tend to have higher star formation efficiency than non-barred centers due to a generally higher velocity dispersion, but such trend is easily obscured using other prescriptions. Contrary to the centers, the disk regions show consistent distribution of velocity dispersion and depletion time between barred and non-barred galaxies, regardless of which prescription is used.

The correlation of $\alpha_{\rm CO}$ with ~ 100 -pc scale velocity dispersion with only a $\sigma \sim 0.1$ dex scatter (see Section 3.1), contrary to ~ 0.3 dex or larger scatter using Z- and/or Σ_{star} -based prescriptions, shows that velocity dispersion is an excellent observational tracer for $\alpha_{\rm CO}$ variations in star-forming galaxies. The rationale behind such a strong relation may be that Δv directly traces the optical depth changes that are the dominant effect responsible for altering $\alpha_{\rm CO}$ across these galaxies, as it has been shown that opacity variation is the primary driver of $\alpha_{\rm CO}$ in various galaxy centers (Israel 2020; Teng et al. 2022, 2023). However, effects of CO-dark gas and CO excitation can also be important to explain $\alpha_{\rm CO}$ variations across the galaxy disks, which have therefore motivated previous $\alpha_{\rm CO}$ prescriptions based on metallicity and/or CO integrated intensity (e.g., Narayanan et al. 2012; Hunt et al. 2015; Amorín et al. 2016; Accurso et al. 2017; Gong et al. 2020).

As discussed in Section 3.2, the correlation of $\alpha_{\rm CO}$ with metallicity (Z') is indirectly included in the dependence with Δv because both Z' and Δv vary with galactocentric radius and are thus correlated. Furthermore, statistical studies on molecular cloud properties have shown that velocity dispersion also correlates well with molecular gas surface density and the CO integrated intensity across galaxy disks (Heyer et al. 2009; Sun et al. 2020b, 2022; Rosolowsky et al. 2021). Therefore, it is likely that our Δv -based prescription contains opacity variations and metallicity gradients as well as the physics of the $\alpha_{\rm CO}$ - $I_{\rm CO}$ correlation suggested by simulation studies (Narayanan et al. 2012; Gong et al. 2020; Hu et al. 2022). This means that the proposed prescription (Equation 2) may incorporate more than one piece of physics into a single scaling relation, which could explain why the trend holds across different galactic environments.

In Sections 3.3 and 3.4, our prescription (based on the dust $\alpha_{\rm CO}$ measurements) suggests lower $\alpha_{\rm CO}$ in barred galaxy centers that lead to higher SFE than non-barred centers and the disks. This low $\alpha_{\rm CO}$ and high SFE in barred centers imply that the amount of molecular gas can be overestimated by previous studies due to inaccurate $\alpha_{\rm CO}$ or the assumption of a constant SFE. By comparing the derived $\Sigma_{\rm mol}$ under different $\alpha_{\rm CO}$ assumptions for all galaxies in Table 1, we find that the median $\Sigma_{\rm mol}$ of barred centers is 3 times higher than that of non-barred centers if using a MW-like $\alpha_{\rm CO}$. On the other hand, our $\alpha_{\rm CO}$ prescription results in only 1.3 times higher $\Sigma_{\rm mol}$ in barred centers. Therefore, it is likely that the enhanced SFE is a more important factor causing high SFR observed in barred galaxy cen-

ters, compared to an increased amount of molecular gas driven inwards by bars.

Recent studies using $\alpha_{\rm CO}$ prescriptions from Narayanan et al. (2012) or B13 also show that barred galaxies tend to have higher central gas concentration than non-barred galaxies, although the degree of concentration is not as significant as using a constant $\alpha_{\rm CO}$ (Sakamoto et al. 1999; Sheth et al. 2005; Schinnerer et al. 2006; Kuno et al. 2007). Such accumulation of gas towards the centers can increase SFR in barred centers, and it is consistent with the theoretical expectation that non-axisymmetric gravitational potential from bars can induce gas inflows and transport more gas into galaxy centers (e.g., Wada & Habe 1995; Regan & Teuben 2004; Kim et al. 2012; Tress et al. 2020b). Bars thus influence the secular evolution of galaxies by redistributing molecular gas mass and angular momentum (see review by Kormendy & Kennicutt 2004).

Studies have also shown that if $\alpha_{\rm CO}$ changes were treated properly, starbursts in galaxy centers and variations of SFRs across nearby galaxies are primarily driven by higher SFE rather than increased molecular gas fraction (Leroy et al. 2013; Ellison et al. 2020a,b; den Brok et al. 2023). This is contrary to studies using constant, Z-based, or Σ_{star} -based α_{CO} , which resulted in similar SFE between barred and non-barred galaxies (e.g., Saintonge et al. 2012; Querejeta et al. 2021, see also Section 3.4). With our proposed $\alpha_{\rm CO}$ prescription, we find enhanced SFE in barred centers, which could originate from variations in molecular gas distribution, density structure, or dynamical effects of turbulence and shocks powered by stellar feedback (e.g., Kainulainen et al. 2009; Renaud et al. 2012; Liu et al. 2023). However, these factors driving SFE variations are the same ones that can alter CO emissivity and $\alpha_{\rm CO}$. Therefore, only with accurate $\alpha_{\rm CO}$ values can we disentangle SFE from $\alpha_{\rm CO}$ and unravel the physical drivers of SFR. Using the latest and best possible measurements of $\alpha_{\rm CO}$ and molecular gas properties across a sample of nearby galaxies, our work lays a foundation for benchmarking $\alpha_{\rm CO}$ calibration in star-forming galaxies (including starbursting galaxy centers) and allows for further investigation on SFE, SFE per cloud free-fall time, or other related properties that can improve our knowledge of galaxy evolution.

5. CONCLUSIONS

We construct a new $\alpha_{\rm CO}$ prescription applicable to star-forming galaxies, where CO emissivity variations are critical in altering $\alpha_{\rm CO}$. The prescription is a major step towards precise calibration of $\alpha_{\rm CO}$ across galaxies, and it reveals unprecedented trends in star formation

properties which may have been obscured by previous $\alpha_{\rm CO}$ prescriptions. Our key results are summarized as follows:

- 1. The strong anti-correlation between measured $\alpha_{\rm CO}$ and CO velocity dispersion (Δv) at ~ 100 pc scales shows that Δv is useful for predicting $\alpha_{\rm CO}$, and it suggests that CO opacity altered by Δv changes or other correlated properties of the molecular gas across the entire galaxies are primary drivers of $\alpha_{\rm CO}$ in star-forming galaxies.
- 2. The proposed α_{CO} prescription (Equation 2) is applicable to regions with metallicity above 0.6 Z_{\odot} and $\langle \Delta v \rangle_{150 \mathrm{pc}} \gtrsim 3 \mathrm{~km~s^{-1}}$. The expected scatter in $\alpha_{\rm CO}$ is $\sigma \sim 0.1$ dex, which is a substantial improvement over existing $\alpha_{\rm CO}$ prescriptions. Our Δv -based prescription has the advantage of connecting directly to the physical causes of $\alpha_{\rm CO}$ change (e.g., CO opacity) as well as requiring only the CO observations which is most relevant to tracing molecular gas.
- 3. With the measured $\alpha_{\rm CO}$, we find distinctly shorter molecular gas depletion time $(t_{\rm dep})$ in barred galaxy centers than non-barred galaxy centers, as well as a generally shorter $t_{\rm dep}$ in galaxy centers than the disks. In contrast, assuming a constant MW $\alpha_{\rm CO}$ results in $t_{\rm dep} \sim 3$ Gyr for all regions, which underestimates the star formation efficiency (SFE) in galaxy centers and also obscures the difference between barred and non-barred galaxies.
- 4. Our prescription reveals short $t_{\rm dep}$ down to 100 Myr in barred galaxy centers, with the median $t_{\rm dep}$ $(0.7^{+0.9}_{-0.4}~{\rm Gyr})$ being 3 times shorter than in nonbarred galaxy centers $(2.1^{+1.5}_{-1.3}~{\rm Gyr})$. However, all other prescriptions (MW, metallicity-based, and B13) show < 0.2 dex difference between the two regions, even if B13 results in an overall shorter $t_{\rm dep}$ for galaxy centers which aligns better with our results. Thus, SFE in barred galaxy centers may be underestimated by a factor of three or more in previous studies due to $\alpha_{\rm CO}$ uncertainties.
- 5. All four prescriptions tested in this work show similar $t_{\rm dep}$ of 2–3 Gyr in the disk regions and nonbarred galaxy centers across the PHANGS sample, which is in good agreement with previous literature (e.g., Leroy et al. 2008; Saintonge et al. 2011; Sun et al. 2023).

This work was carried out as part of the PHANGS collaboration. Y.-H.T. and K.S. acknowledge funding support from NRAO Student Observing Support Grant SOSPADA-012 and from the National Science Foundation (NSF) under grant No. 2108081. Y.-H.T. also acknowledges support by the Ministry of Education of Taiwan through Government Scholarship to Study J.S. acknowledges support by the Natural Sciences and Engineering Research Council of Canada (NSERC) through a Canadian Institute for Theoretical Astrophysics (CITA) National Fellowship. KG is supported by the Australian Research Council through the Discovery Early Career Researcher Award (DECRA) Fellowship (project number DE220100766) funded by the Australian Government. KG is supported by the Australian Research Council Centre of Excellence for All Sky Astrophysics in 3 Dimensions (ASTRO 3D), through project number CE170100013. HAP acknowledges support by the National Science and Technology Council of Taiwan under grant 110-2112-M-032-020-MY3. MB gratefully acknowledges support from the ANID BASAL project FB210003 and from the FONDE-CYT regular grant 1211000. IC thanks the National Science and Technology Council for support through grant 111-2112-M-001-038-MY3, and the Academia Sinica for Investigator Award AS-IA-109-M02. JC acknowledges funding from the Belgian Science Policy Office (BEL-SPO) through the PRODEX project "JWST/MIRI Science exploitation" (C4000142239). MC gratefully acknowledges funding from the DFG through an Emmy Noether Research Group (grant number CH2137/1-COOL Research DAO is a Decentralized Autonomous Organization supporting research in astrophysics aimed at uncovering our cosmic origins. SKS acknowledges financial support from the German Research Foundation (DFG) via Sino-German research grant SCHI 536/11-1. AU acknowledges support from the Spanish grant PID2019-108765GB-I00, funded by MCIN/AEI/10.13039/501100011033.

This paper makes use of the following ALMA data: ADS/JAO.ALMA#2012.1.00650.S,

ADS/JAO.ALMA#2013.1.00803.S,

ADS/JAO.ALMA#2013.1.01161.S,

ADS/JAO.ALMA#2015.1.00925.S,

ADS/JAO.ALMA#2015.1.00956.S,

ADS/JAO.ALMA#2017.1.00392.S.

ADS/JAO.ALMA#2017.1.00766.S,

ADS/JAO.ALMA#2017.1.00886.L, ADS/JAO.ALMA#2018.1.01321.S,

ADS/JAO.ALMA#2018.1.01651.S.

ADS/JAO.ALMA#2018.A.00062.S.

ALMA is a partnership of ESO (representing its member states), NSF (USA), and NINS (Japan), together with NRC (Canada), NSC and ASIAA (Taiwan), and KASI (Republic of Korea), in cooperation with the Republic of Chile. The Joint ALMA Observatory is operated by ESO, AUI/NRAO, and NAOJ. The National Radio Astronomy Observatory is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc. This work is based in part on data products from the Wide-field Infrared Survey Explorer (Wright et al. 2010), which is a joint project of the University of California, Los Angeles, and

the Jet Propulsion Laboratory/California Institute of Technology, funded by the National Aeronautics and Space Administration. We acknowledge the usage of the SAO/NASA Astrophysics Data System.

Facilities: ALMA, WISE, Herschel, VLA

Software: matplotlib (Hunter 2007), numpy (Harris et al. 2020), scipy (Virtanen et al. 2020), astropy (Astropy Collaboration et al. 2013, 2018), ipython (Pérez & Granger 2007)

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